

# Lecture 5

## High Power Microwave Sources

### EEC746

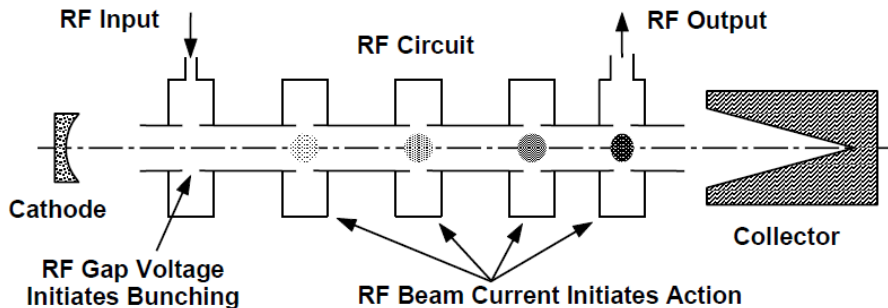
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## 1 Induced Current

- Total Current, Continuity of Total Current
- Schockley-Ramo Theorem

# Beam Gap Interaction



## 1 Induced Current

- Total Current, Continuity of Total Current
- Shockley-Ramo Theorem

# Definition of Total Current

$$\nabla \times \mathbf{H} = \mathbf{J}_c + \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} = \mathbf{J}_c + \mathbf{J}_d$$

- $\mathbf{J}_c$  is the convection current (due to motion of charges).
- $\mathbf{J}_d$  is the displacement current (due to time variation of the electric field).

## Definition

Total current density  $\mathbf{J}_t$  is defined as,

$$\mathbf{J}_t \equiv \mathbf{J}_c + \mathbf{J}_d$$

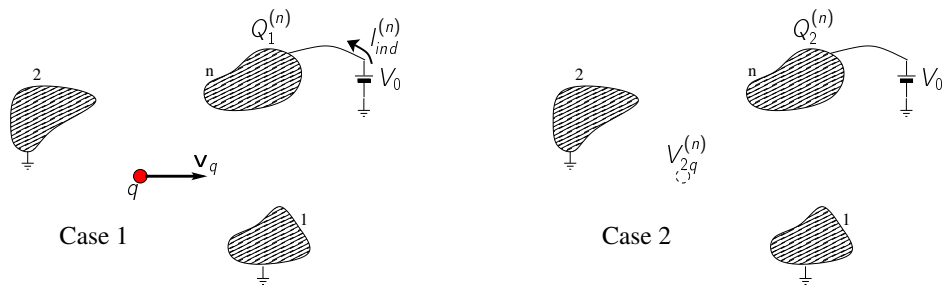
## Continuity of total current

$$\nabla \cdot \mathbf{J}_t = 0$$

## 1 Induced Current

- Total Current, Continuity of Total Current
- Schockley-Ramo Theorem

# Schockley-Ramo Theorem



Green's Theorem for the two cases potential functions  $V_1$  and  $V_2$ ,

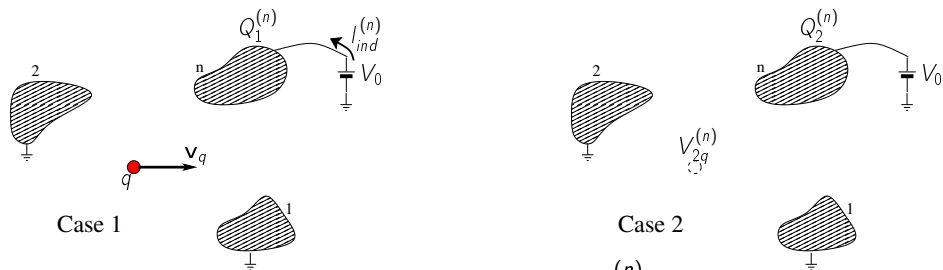
$$\int_V (V_1 \nabla^2 V_2 - V_2 \nabla^2 V_1) d\tau = \oint_S (V_1 \nabla V_2 - V_2 \nabla V_1) \cdot \mathbf{n} da$$

$$\nabla^2 V_1 = -\frac{q}{\epsilon_0} \delta^3(\mathbf{r} - \mathbf{r}_q), \quad \nabla^2 V_2 = 0$$

$$qV_{2q}^{(n)} = Q_1^{(n)}V_0 - Q_2^{(n)}V_0$$

$$Q_{ind}^{(n)} = Q_2^{(n)} - Q_1^{(n)} = -q \frac{V_{2q}^{(n)}}{V_0}$$

# Schockley-Ramo Theorem



$$Q_{ind}^{(n)} = Q_2^{(n)} - Q_1^{(n)} = -q \frac{V_{2q}^{(n)}}{V_0}$$
$$I_{ind}^{(n)} = \frac{dQ_{ind}^{(n)}}{dt} = -\frac{q}{V_0} \mathbf{v}_q \cdot \nabla V_{2q}^{(n)} = \frac{1}{V_0} q \mathbf{v}_q \cdot \mathbf{E}_q$$

For a group of charges,

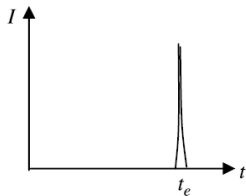
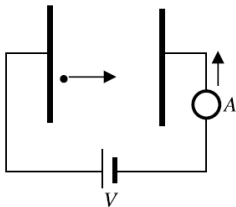
$$I_{ind}^{(n)} = \frac{1}{V_0} \sum_q q \mathbf{v}_q \cdot \mathbf{E}_q = \frac{1}{V_0} \int_V \mathbf{J}_c \cdot \mathbf{E} d\tau$$

$$I_{ind}^{(n)} = \frac{1}{V_0} \int_V \mathbf{J}_c \cdot \mathbf{E} d\tau$$

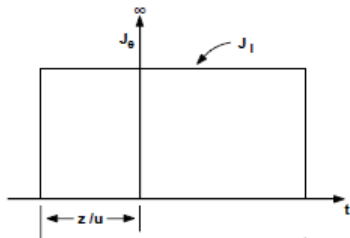


# Gridded (Planar) Gaps

Electron Bunch with Uniform Velocity



$$l_i = \frac{1}{V} \int_V \mathbf{J}_c \cdot \mathbf{E} d\tau, \quad l_i = \frac{1}{-El} (-\sigma) uEA = \frac{\sigma}{l} uA$$

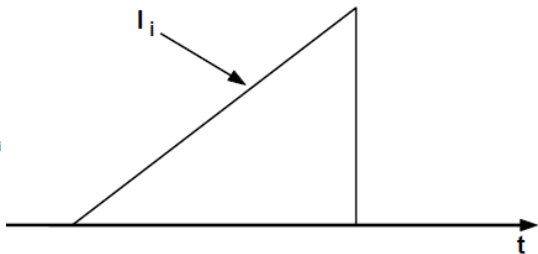
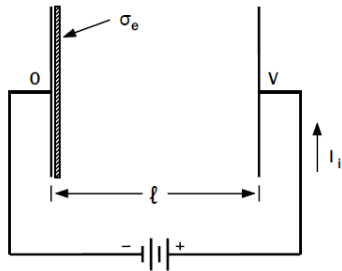


# Gridded (Planar) Gaps

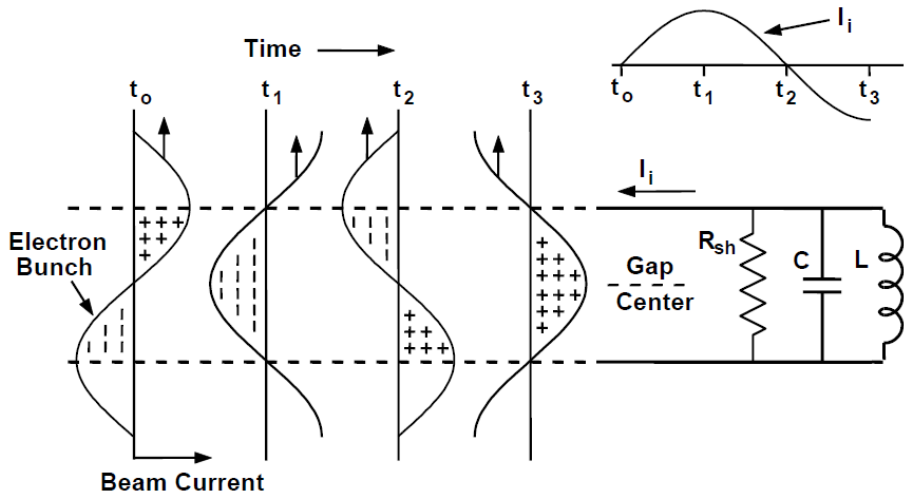
Change in Velocity is Considered

$$u = \eta \frac{V}{\ell} t,$$

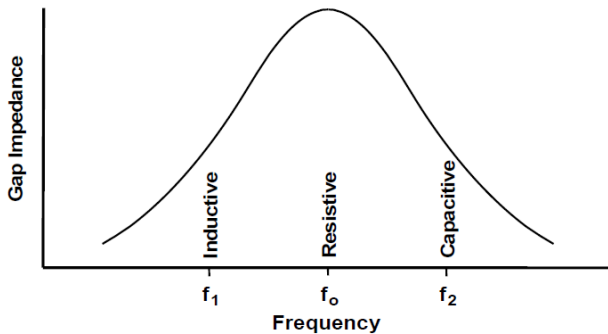
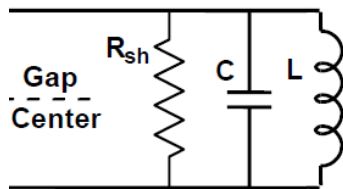
$$I_i = \frac{1}{V} \int \mathbf{J}_c \cdot \mathbf{E} d\tau = \frac{1}{V} \eta (-\sigma) \eta \frac{V}{\ell} t \times \left( -\frac{V}{\ell} \right) A = \sigma \eta \frac{V}{\ell^2} A t$$



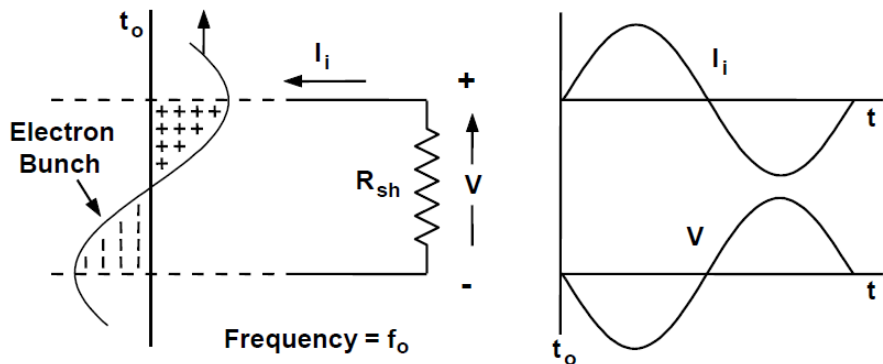
# Current Induced in Cavity Circuit



# Impedance Presented to the Beam as a Function of Frequency



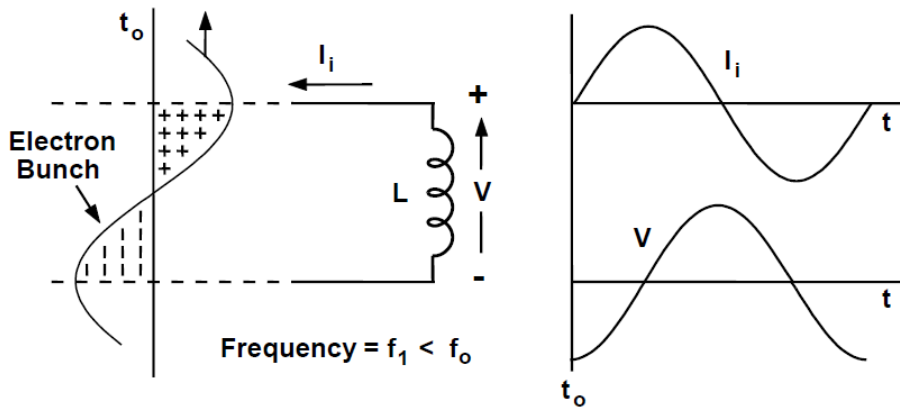
# Voltage Induced with Resistive Loading



$$I_i = I_{\max} \sin \omega t, \quad V = -I_i R = -I_{\max} R \sin \omega t$$

- The induced voltage slightly decelerates the electron bunches.
- A new signal is generated in the beam which is  $90^\circ$  with the current  $I_i$ .

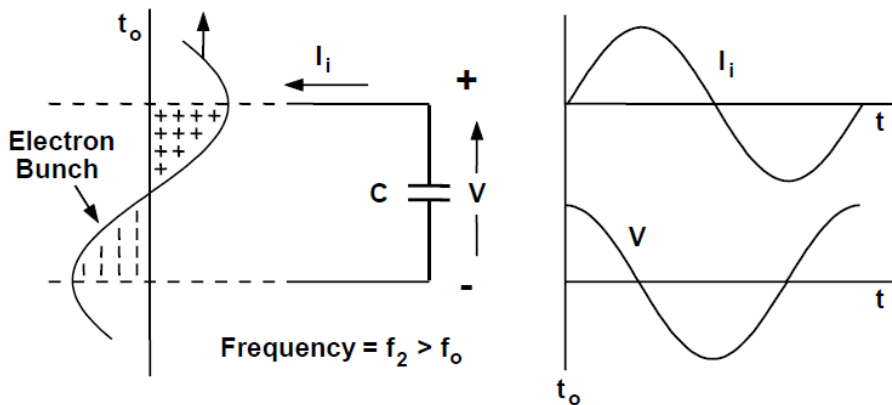
# Voltage Induced with Inductive Loading



$$I_i = I_{\max} \sin \omega t, \quad V = \Im \{ -j\omega L I_{\max} e^{j\omega t} \} = -\omega L I_{\max} \cos \omega t$$

The induced voltage decelerates the electrons in and near the leading edge and accelerates those in and near the trailing edge, hence enhances the beam bunching.

# Voltage Induced with Capacitive Loading



$$I_i = I_{\max} \sin \omega t, \quad V = \Im \left\{ \frac{-1}{j\omega C} I_{\max} e^{j\omega t} \right\} = \frac{I_{\max}}{\omega C} \cos \omega t$$

The induced voltage accelerates electrons away from the leading and trailing edges of a bunch, which tends to destroy the bunch (debunching).