

Lecture 3

High Power Microwave Sources

EEC746

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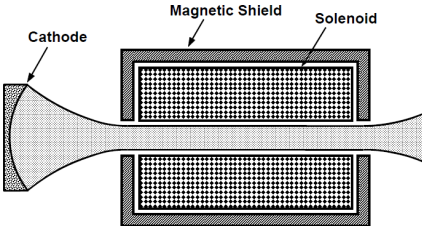
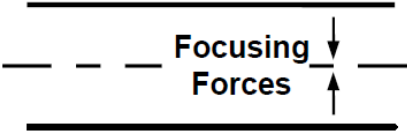
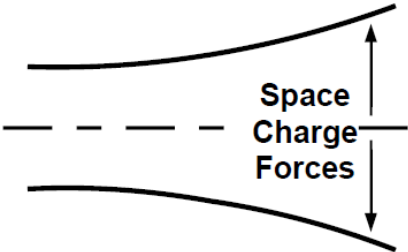
1 Overview of Uniform Field Focusing

- Brillouin Flow
- Scalloping
- Confined (Immersed) Flow

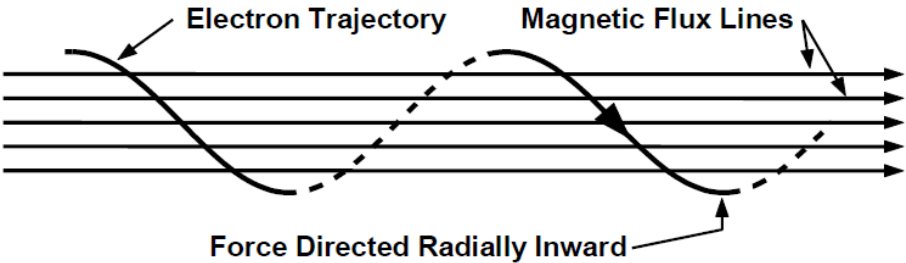
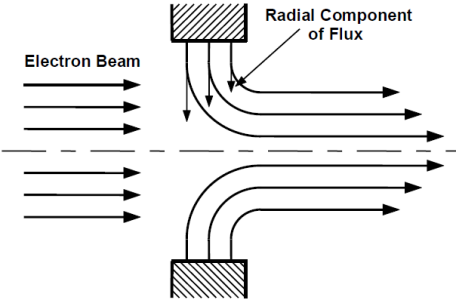
2 Uniform-Field Focusing and Laminar Flow

- The Beam Equation
- Confined (Immersed) Flow

Electron Dynamics

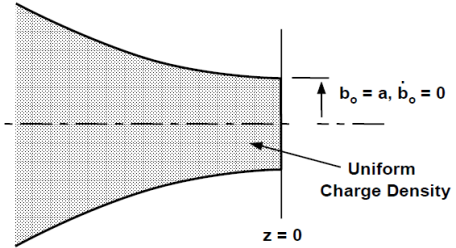
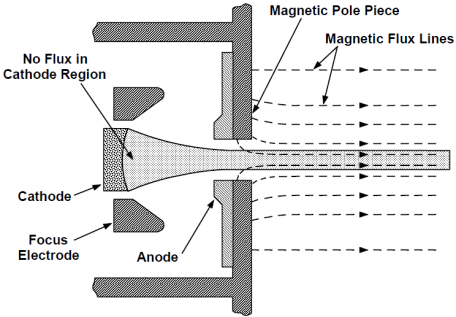


Electron Dynamics



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Brillouin Flow



Brillouin Flow

$$F_r = m(\ddot{r} - r\dot{\theta}^2) = -eE_r - er\dot{\theta}B_z$$

On the beam envelope $r = b$,

$$E_r = \frac{-I}{2\pi\epsilon_0 u_0 b}$$

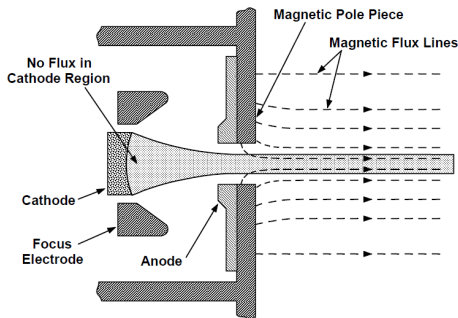
$$\ddot{b} - b\dot{\theta}^2 = \frac{\eta I}{2\pi\epsilon_0 u_0 b} - \eta b\dot{\theta}B_z$$

According to Busch's theorem, $\dot{\theta} \approx \frac{\eta}{2} \left(B_z - \frac{r_0^2}{r^2} B_{z0} \right)$

As in the cathode region $B_{z0} = 0$, $\dot{\theta} = \frac{\eta}{2} B_z = \omega_L \equiv$ Larmer frequency

$$\ddot{b} + \omega_L^2 b = \frac{\eta I}{2\pi\epsilon_0 u_0 b}$$

For Brillouin flow the equilibrium beam radius ($\ddot{b} = 0$), $a = \frac{1}{B_z} \sqrt{\frac{2I}{\pi\epsilon_0 \eta u_0 b}}$.



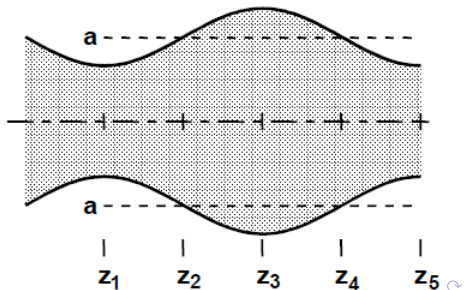
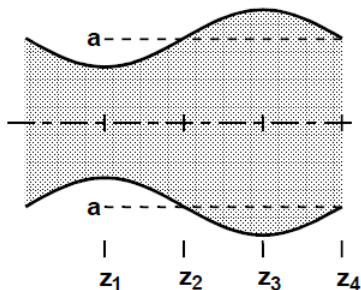
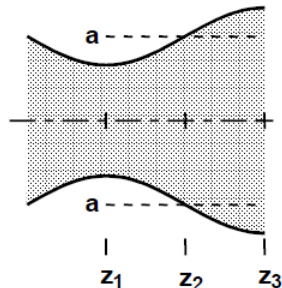
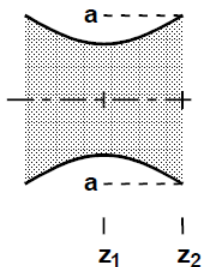
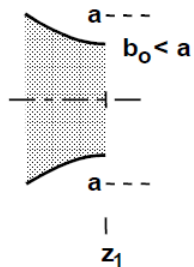
1 Overview of Uniform Field Focusing

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- **Scalloping**
- Confined (Immersed) Flow

2 Uniform-Field Focusing and Laminar Flow

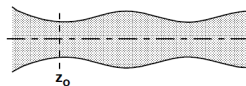
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Scalloping

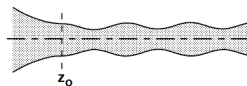


Brillouin Flow Problems

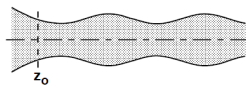
- 1 Flux density is too small.
- 2 Flux density is too large.
- 3 Beam is converging at the entrance.
- 4 Beam is diverging at the entrance.
- 5 Beam axis is offset from the magnetic field axis.
- 6 Beam axis is tilted with respect to the magnetic field axis.



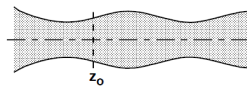
(a) $B < B_B$



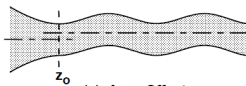
(b) $B > B_B$



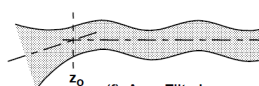
(c) Beam Converging



(d) Beam Diverging



(e) Axes Offset



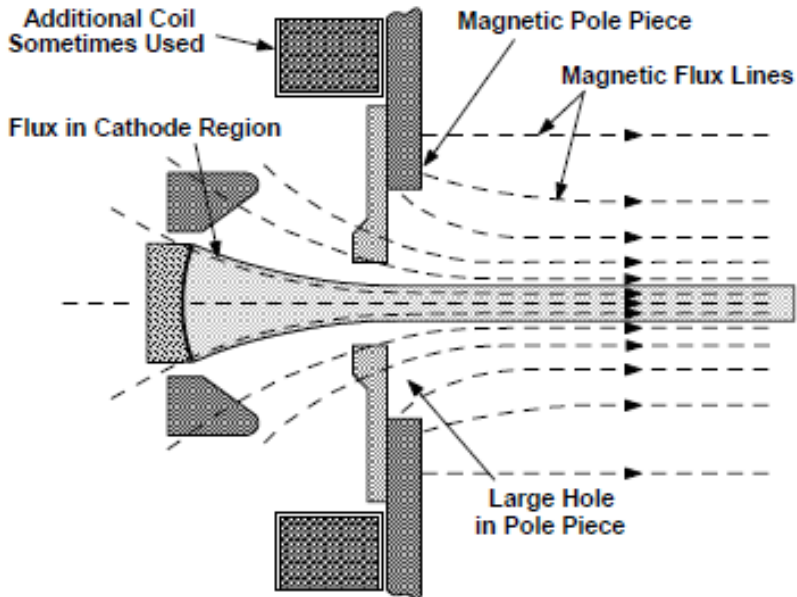
(f) Axes Tilted

Brillouin Flow Pros and Cons

- Magnetic field required is the lower for any other focusing system.
- Beam is extremely sensitive to misalignment and perturbation.
- Brillouin flow can be very nearly achieved under laboratory conditions. In practice, the magnetic focusing field that is used, even when the cathode is shielded, is greater than the Brillouin value. Reasons for this include transverse electron velocities (from thermal and other effects in the gun) and the increase in beam size that results from RF modulation.

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Confined (Immersed) Flow



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The Beam Equation

$$\ddot{r} - r\dot{\theta}^2 = -\eta \left(E_r + Br\dot{\theta} \right),$$

where $\dot{\theta}$ is given by Busch's Theorem,

$$\dot{\theta} = \frac{\eta}{2} \left(B - B_c \frac{r_c^2}{r^2} \right)$$

The electric field on the beam boundary $r = b$,

$$E_r(b) = \frac{-I}{2\pi b \epsilon_0 u_0},$$

The envelope equation is given by,

$$\ddot{b} - b \frac{\eta^2}{4} \left(B - B_c \frac{b_c^2}{b^2} \right)^2 = -\eta \left[\frac{-I}{2\pi b \epsilon_0 u_0} + Bb \frac{\eta}{2} \left(B - B_c \frac{b_c^2}{b^2} \right) \right]$$
$$\ddot{b} + b \omega_L^2 \left(1 - \frac{B_c^2 b_c^4}{B^2 b^4} \right) - \frac{\eta I}{2\pi b \epsilon_0 u_0} = 0$$

The Beam Equation

$$\ddot{b} + b\omega_L^2 \left(1 - \frac{B_c^2 b_c^4}{B^2 b^4} \right) - \frac{\eta l}{2\pi b \epsilon_0 u_0} = 0$$

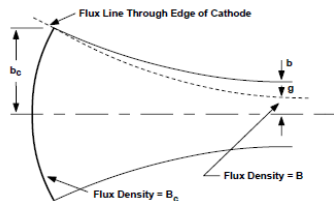
$$\pi b_c^2 B_c = \pi g^2 B \quad \Rightarrow \quad \frac{B_c b_c^2}{B b^2} = \frac{g^2}{b^2}$$

Using the Brillouin radius $a = \frac{1}{\omega_L} \left(\frac{\eta l}{2\pi \epsilon_0 u_0} \right)^{1/2}$,

$$\ddot{b} + \omega_L^2 \left[b \left(1 - \frac{g^4}{b^4} \right) - \frac{a^2}{b} \right] = 0$$

Normalizing the radii, $R = \frac{b}{a}$, and $R_g = \frac{g}{a}$,

$$\ddot{R} + \omega_L^2 \left[R \left(1 - \frac{R_g^4}{R^4} \right) - \frac{1}{R} \right] = 0$$



The Beam Equation

$$\ddot{R} + \omega_L^2 \left[R \left(1 - \frac{R_g^4}{R^4} \right) - \frac{1}{R} \right] = 0$$

Equilibrium radius R_e corresponds to $\ddot{R}_e = 0$, and is given by the following equation,

$$R_e \left(1 - \frac{R_g^4}{R_e^4} \right) - \frac{1}{R_e} = 0 \quad \Rightarrow \quad R_e = \left[\frac{1}{2} + \frac{1}{2} \sqrt{1 + 4R_g^4} \right]^{1/2}$$

The Beam Equation

$$\ddot{R} + \omega_L^2 \left[R \left(1 - \frac{R_g^4}{R^4} \right) - \frac{1}{R} \right] = 0$$

So the solution near the equilibrium $R = R_e(1 + \delta)$ can be linearized as follows,

$$\ddot{\delta} + 2\Omega\omega_L^2\delta = 0, \quad \text{where } \Omega = 2 - \frac{1}{R_e^2}$$

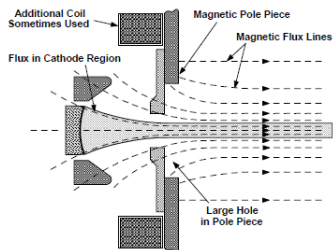
$$\delta = A \cos(\sqrt{2\Omega}\omega_L t) + A \sin(\sqrt{2\Omega}\omega_L t)$$

$$R = R_e \left[1 + A \cos\left(\sqrt{2\Omega}\frac{\omega_L}{u_0} z\right) + A \sin\left(\sqrt{2\Omega}\frac{\omega_L}{u_0} z\right) \right]$$

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Confined (Immersed) Flow



Absence of flux in cathode region

Presence of flux in cathode region

$$-b_0 \dot{\theta}_B^2 = -\eta (E_r + B_B b_0 \dot{\theta}_B),$$

$$\dot{\theta}_B = \frac{\eta}{2} B_B$$

$$-b_0 \dot{\theta}^2 = -\eta (E_r + B b_0 \dot{\theta})$$

$$\dot{\theta} = \frac{\eta}{2} \left(B - B_c \frac{b_c^2}{b_0^2} \right)$$

$$b_0 \dot{\theta}_B (\eta B_B - \dot{\theta}_B) = b_0 \dot{\theta} (\eta B - \dot{\theta})$$

$$B^2 = B_B^2 + B_c^2 \frac{b_c^2}{b_0^2}$$

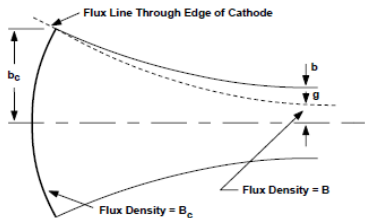
Confined (Immersed) Flow

$$B^2 = B_B^2 + B_c^2 \frac{b_c^2}{b_0^2}$$

Definition

The confinement factor m is defined as,

$$B = mB_B$$



$$b_c^2 B_c = g^2 B = g^2 m B_B, \quad \rightarrow g = \left(1 - \frac{1}{m^2}\right)^{1/4} b_0, \quad \rightarrow R_g = \left(1 - \frac{1}{m^2}\right)^{1/4} R_e$$

$$\therefore R_e \left(1 - \frac{R_g^4}{R_e^4}\right) - \frac{1}{R_e} = 0, \quad \therefore R_e = m,$$

$$R_g = [m^2 (m^2 - 1)]^{1/4}$$